

RESEARCH ARTICLE

Deterministic and replaceable transfer of silver flakes for microcavities

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Supporting Information

S1. Optical image of continuous silver flakes



Fig. S1 Optical image of large-scale continuous silver flakes transferred on the polydimethylsiloxane (PDMS) substrate.

S2. Transferability of silver mirrors of different thicknesses

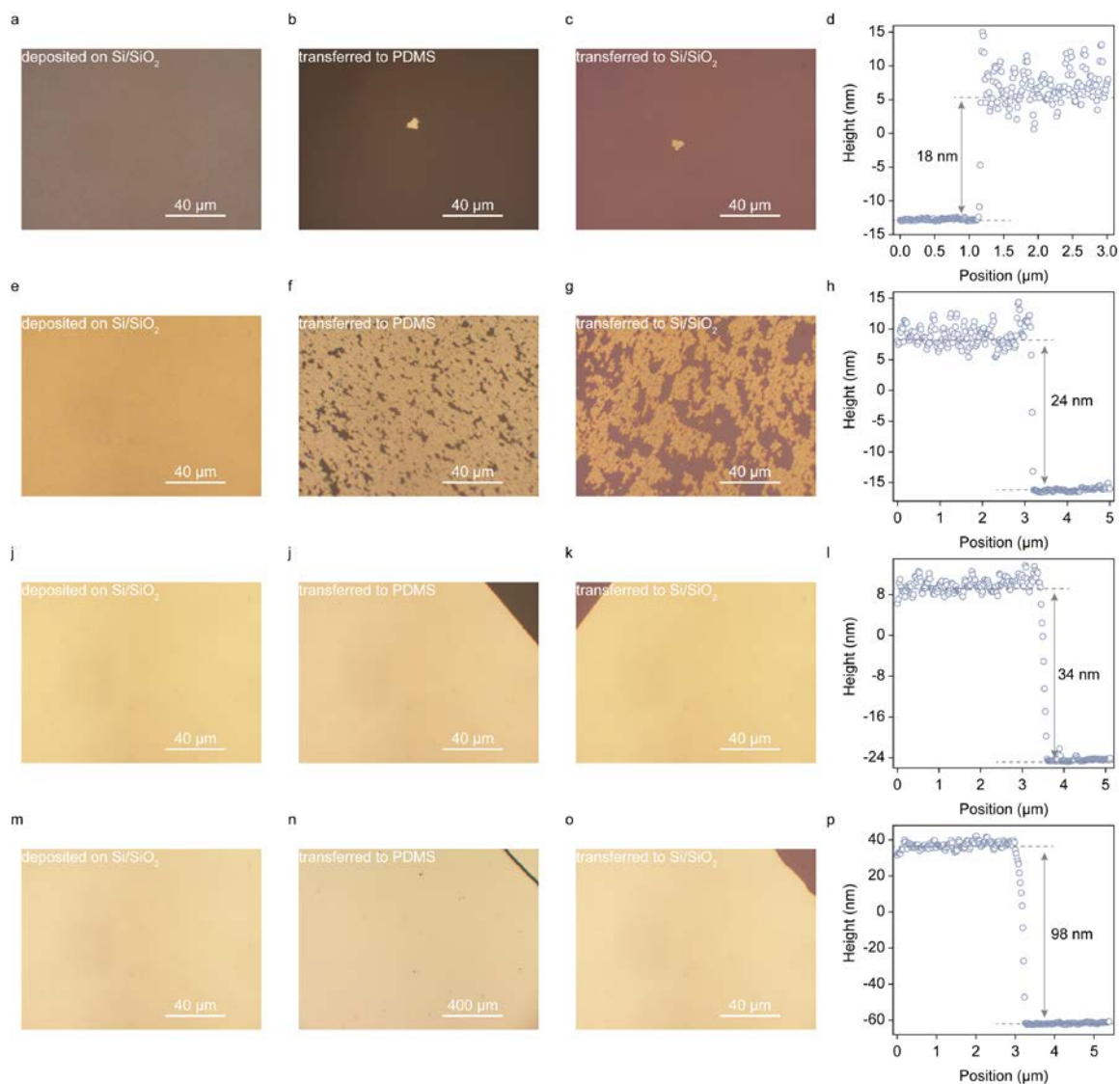


Fig. S2 Deterministic transfer and thickness characterization of silver flakes with different thicknesses. **(a-c, e-g, i-k, and m-o)** Optical microscope images of 18 nm, 24 nm, 34 nm and 98 nm thick silver films deposited on Si/SiO₂ substrates, transferred to PDMS substrates, and transferred to target Si/SiO₂ substrates, respectively. **(d, h, i, p)** The corresponding height line profiles of the transferred silver flakes on target Si/SiO₂ substrates.

S3. Deterministic transfer of gold flakes

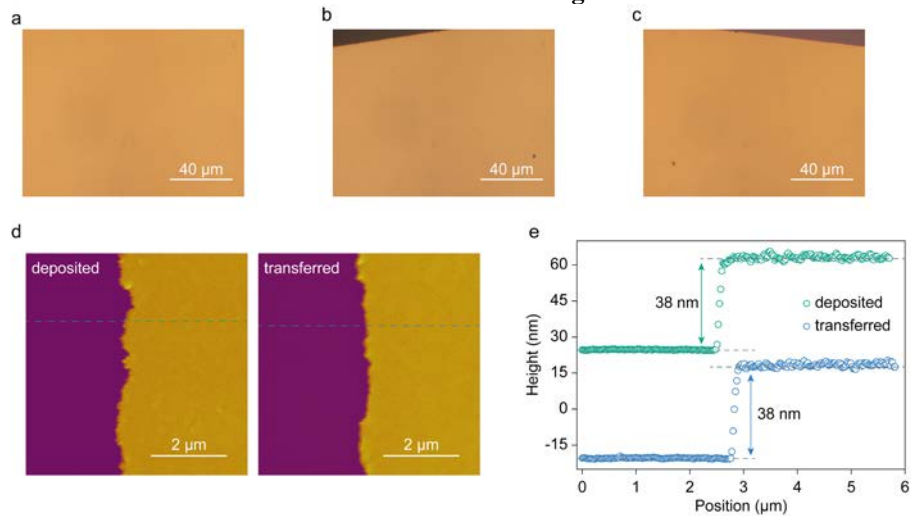


Fig. S3 Deterministic transfer and thickness characterization of gold flakes. **(a-c)** Optical microscope images of gold film deposited on Si/SiO₂ substrate (a), peeled off by PDMS substrate (b), and transferred onto another Si/SiO₂ substrate (c), respectively. **(d-e)** Atomic force microscope height images and the corresponding height line profiles (along the dashed lines marked in (d)) of the deposited and transferred gold flakes.

S4. Roughness characterization of silver flakes

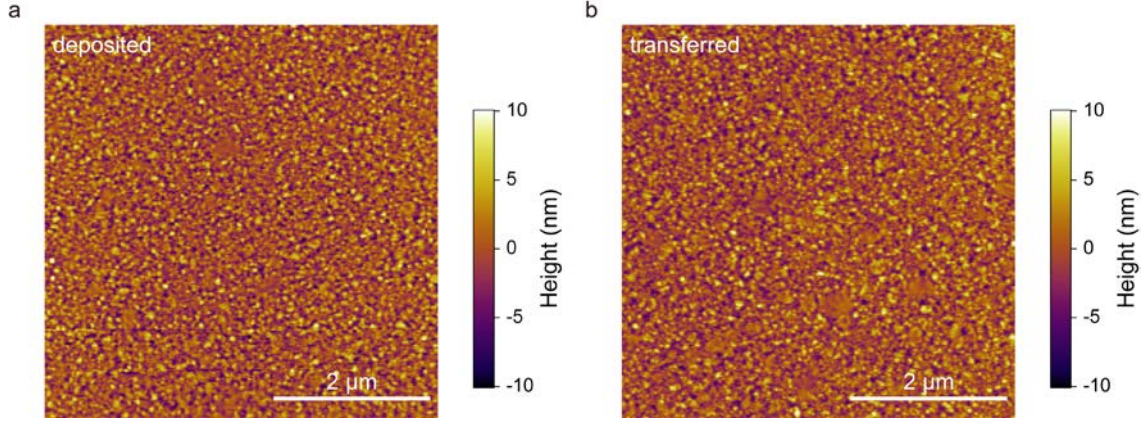


Fig. S4 Surface roughness characterizations of the deposited silver flake (a) and the transferred silver flake (b). Within the area of $25.3 \mu\text{m}^2$, the surface roughness of the silver flake deposited on the Si/SiO₂ substrate is $R_a = 2.047 \text{ nm}$, while the surface roughness of the silver flake transferred to the Si/SiO₂ substrate is $R_a = 1.975 \text{ nm}$.

Note S1. Dispersion fitting

Traditionally, strong coupling behavior can be described by a simplified coupled oscillator model that ignores linewidth:

$$\begin{bmatrix} E_c & V \\ V & E_{ex} \end{bmatrix} \begin{bmatrix} a \\ b \end{bmatrix} = E \begin{bmatrix} a \\ b \end{bmatrix}, \quad (1)$$

where E_c is the cavity mode (E_{ca} and E_{cb} correspond to cavity modes along the a - and b -axis, respectively), E_{ex} is the exciton energy, E is the eigenvalue corresponding to the polariton mode energy, a and b define the eigenvectors of the polariton modes, and V is the coupling strength between the exciton state and photonic mode. Owing to the multilayer structure of our device, we can derive the cavity mode as

$$E_c(\theta) = \frac{mhc}{\sum_i \frac{2n_i L_i}{\sqrt{1 - \left(\frac{\sin\theta}{n_i}\right)^2}}}, \quad (2)$$

where m is the mode order, θ is the emission angle relative to the z -axis, n_i and L_i are the refractive index and thickness of the i th layer, respectively.

To avoid fitting errors, we obtain a series of material parameters from our experiments or references to reduce the number of free-tuning parameters. The cavity modes are calculated from the fixed thickness and refractive index of each layer in the cavity, while the exciton energies are directly extracted from their absorption spectra at different temperatures.

In this way, for the microcavity in Fig. 2(d), the degenerate exciton energy extracted from the room temperature absorption spectrum is 2.3574 eV, and the best fit yields a cavity mode of $E_{ca} = 2.3176 \text{ eV}$ along the a -axis and a cavity mode of $E_{cb} = 2.3297 \text{ eV}$ along the b -axis. In addition, we obtain the Rabi splitting energy, $\hbar\Omega$, between the lower polariton (LP) branch and the upper polariton (UP) branch along the a -axis and b -axis, which are 107.94 meV and 96.23 meV, respectively.

Similarly, for the microcavities in Figs. 3(h) and (i), the degenerate exciton energy is maintained at 2.3608 eV. As for the microcavity in Fig. 3(h), the best fit yields a cavity mode of $E_{ca} = 2.3978 \text{ eV}$ along the a -axis and a cavity mode of $E_{cb} = 2.4061 \text{ eV}$ along the b -axis. In addition, we obtain the Rabi splitting energy, $\hbar\Omega$, between the LP branch and the UP branch along the a -axis and b -axis, which are 53.45 meV and 42.19 meV, respectively. After the subsequent transfer steps, for the microcavity in Fig. 3(i), the best fit yields a cavity mode of $E_{ca} = 2.3561 \text{ eV}$ along the a -axis and a cavity mode of $E_{cb} = 2.3641 \text{ eV}$ along the b -axis. In addition, we obtain the Rabi splitting energy, $\hbar\Omega$, between the LP branch and the UP branch along the a -axis and b -axis, which are 68.77 meV and 60.83 meV, respectively.

S5. Fitting of exciton-polariton emission

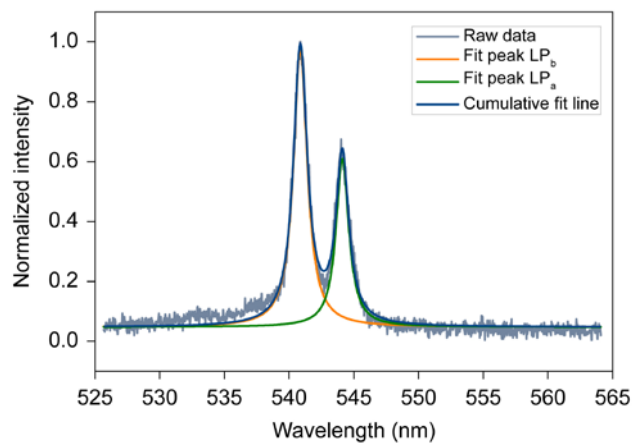


Fig. S5 Photoluminescence spectrum of exciton-polariton emission recorded at $\theta = 0^\circ$ (shadow blue line) and its corresponding fitting lines. The yellow, green and safety blue lines display the fittings of the LP dispersion LP_b, LP_a, and the cumulative peak, respectively. LP_a (LP_b) is centered at 544.1 nm (540.9 nm) with an full width at half maximum of 1.20 nm (1.29 nm), corresponding to a quality factor of ~ 454 (~ 420).

S6. Optical image of the top h-BN layer on PDMS substrate

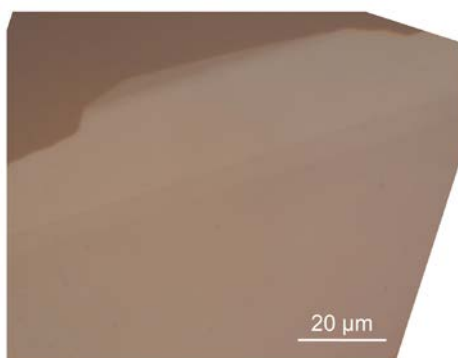


Fig. S6 The optical image of the h-BN layer exfoliated on PDMS substrate, which helps outline its profile in the optical image of the device in Fig. 3(f) in the main text.

S7. Hopfield coefficients

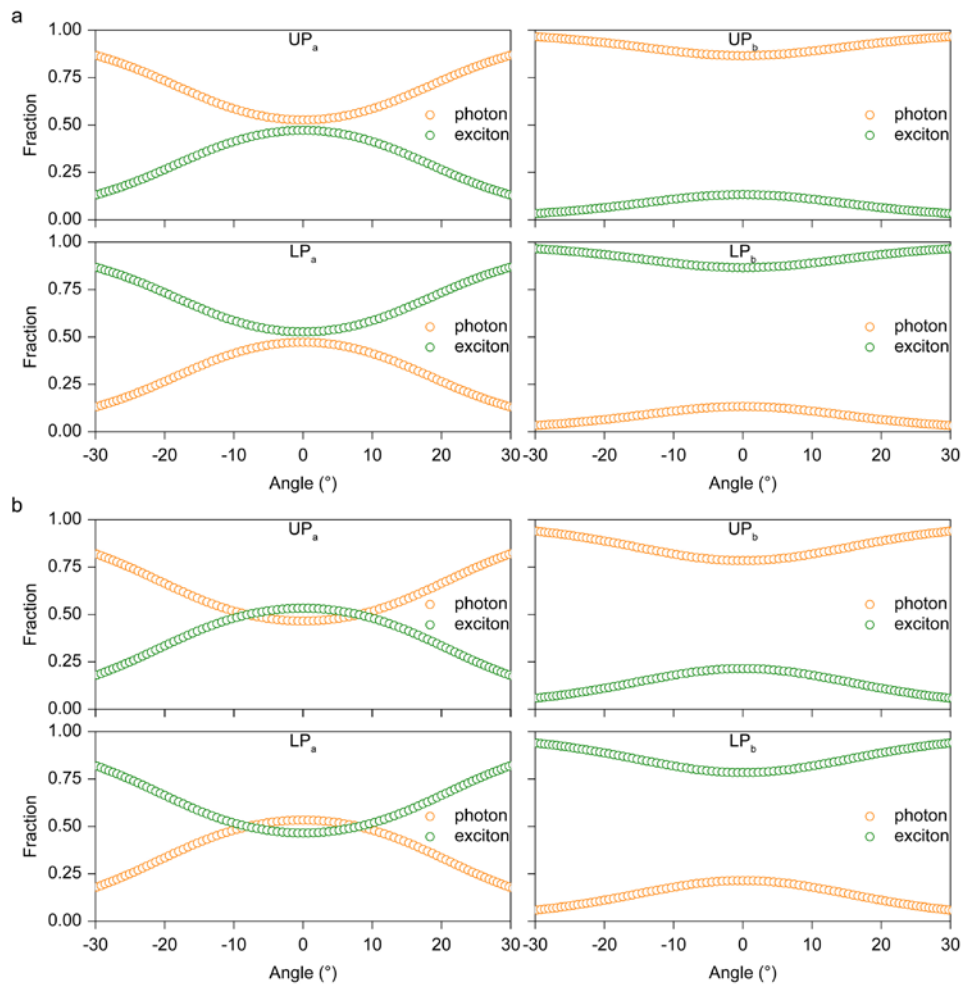


Fig. S7 Fitted Hopfield coefficients. **(a, b)** Fitted Hopfield coefficients of the cavity photons and excitons for the polariton branches with polarizations along the a -axis and b -axis in Fig. 3(h) (a) and Fig. 3(i) (b) in the main text.